## Communication

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# Mapping Chemical Exchange in Proteins with MW > 50 kD 

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Chemical exchange phenomena in nuclear magnetic resonance (NMR) spectroscopy of proteins in solution arise from motions on $\mu \mathrm{s}$ to ms time scales that can be critical for ligand binding, ${ }^{1}$ catalysis, ${ }^{2}$ and allosteric regulation. ${ }^{3}$ Use of chemical exchange to study these biological processes usually has been restricted to proteins with molecular weight (MW) $<25 \mathrm{kD} .{ }^{4}$ The chemical exchange contribution to transverse relaxation, $R_{\mathrm{ex}}$, is independent of protein size and consequently is a smaller fraction of the apparent transverse relaxation rate constant, $R_{2}$, for larger proteins. In addition, NMR spectroscopy of larger proteins is hindered by severe resonance overlap and rapid transverse relaxation. This Communication describes a novel experiment for rapid identification of chemical exchange in ${ }^{2} \mathrm{H} /{ }^{15} \mathrm{~N}$-labeled proteins with MW ranging from 15 to $>50 \mathrm{kD}$. The technique is demonstrated in Escherichia coli ribonuclease H ( $\mathrm{RNaseH)} \mathrm{a} \mathrm{monomer} \mathrm{with} \mathrm{MW}=,17 \mathrm{kD}$, and Saccharomyces cerevisiae triosephosphate isomerase (TIM), a symmetric dimer with total MW $=54 \mathrm{kD}$.

Interference between ${ }^{1} \mathrm{H}^{-15} \mathrm{~N}$ dipole-dipole (DD) and ${ }^{15} \mathrm{~N}$ chemical shift anisotropy (CSA) interactions results in different relaxation rate constants and line width for the two components of the scalar-coupled amide ${ }^{15} \mathrm{~N}$ doublet. The transverse relaxation rate constants, $R_{2}{ }^{\alpha}$ and $R_{2}{ }^{\beta}$, for the narrow (corresponding to the $\mathrm{H}^{\alpha} \mathrm{N}^{-}$ transition) and broad (corresponding to the $\mathrm{H}^{\beta} \mathrm{N}^{-}$transition) doublet components, respectively, are given by: ${ }^{5,6}$

$$
\begin{gather*}
R_{2}^{\alpha}=R_{2}^{0}-\eta_{x y}+R_{1}^{\mathrm{H}} / 2+R_{\mathrm{ex}} \\
R_{2}^{\beta}=R_{2}^{\alpha}+2 \eta_{x y} \tag{1}
\end{gather*}
$$

where $\mathrm{R}_{2}{ }^{0}$ is the auto-relaxation rate constant for in-phase ${ }^{15} \mathrm{~N}$ magnetization from ${ }^{15} \mathrm{~N}$ CSA and ${ }^{1} \mathrm{H}-{ }^{15} \mathrm{~N}$ DD interactions, $\eta_{x y}$ is the rate constant for cross-correlation between ${ }^{15} \mathrm{~N}$ CSA and ${ }^{1} \mathrm{H}-$ ${ }^{15} \mathrm{~N}$ DD interactions, and $R_{1}{ }^{\mathrm{H}}$ is the longitudinal relaxation rate for the ${ }^{1} \mathrm{H}^{\mathrm{N}}$ spin from CSA and DD interactions with remote ${ }^{1} \mathrm{H}$ spins. For ${ }^{2} \mathrm{H} /{ }^{15} \mathrm{~N}$-labeled proteins at high magnetic fields, $R_{2}{ }^{\alpha} \ll R_{2}=$ $R_{2}{ }^{0}+R_{\text {ex }}$ which facilitates the detection of $R_{\text {ex }}$. Although $R_{2}{ }^{0}$ cannot be directly measured for exchanging sites, $\kappa=R_{2}{ }^{0} \eta_{x y}$ is independent of chemical exchange, local motions or global dynamics in large proteins. ${ }^{7,8}$ Thus, $R_{\text {ex }}$ is determined by:

$$
\begin{equation*}
R_{\mathrm{ex}}=R_{2}^{\alpha}-R_{1}^{2 \mathrm{HzNz}} / 2-\eta_{x y}(\kappa-1)+R_{1}^{\mathrm{N} / 2} \tag{2}
\end{equation*}
$$

in which $R_{1}{ }^{2 \mathrm{HzNz}} \approx R_{1}{ }^{\mathrm{H}}+R_{1}{ }^{\mathrm{N}}$ is the relaxation rate constant for longitudinal two-spin order and $R_{1}{ }^{\mathrm{N}}$ is the longitudinal relaxation rate of the ${ }^{15} \mathrm{~N}$ spin. Values of $\kappa$ are derived from the trimmed mean of $1+\left(R_{2}{ }^{\alpha}-R_{1}{ }^{2 \mathrm{HzNz} / 2}\right) / \eta_{x y},{ }^{9}$ or theoretical calculations. ${ }^{7,8}$ $R_{1}{ }^{\mathrm{N}}$ is small and can be neglected in eq 2 for proteins with MW $>$ 30 kD (e.g., $\left\langle R_{1}{ }^{\mathrm{N}}\right\rangle=0.44 \pm 0.04 \mathrm{~s}^{-1}$ in TIM); for smaller proteins, $R_{1}{ }^{\mathrm{N}}$ can be measured using conventional experiments. ${ }^{10}$

[^0]

Figure 1. Pulse sequences for measuring $R_{2}{ }^{\alpha}, R_{2}{ }^{\beta}$, and $R_{1}{ }^{2 H z N z}$ to detect chemical exchange in large proteins. The sequence shown in (a) detects relaxation of the narrow doublet component during the Hahn echo period $2 \tau$ when the proton composite pulse element (open bar) is $\left(90^{\circ}{ }_{x} 90^{\circ}{ }_{y} 90^{\circ}{ }_{-y} 90^{\circ}{ }_{-x}\right.$ ) and detects relaxation of the broad doublet component when this element is $\left(90^{\circ}{ }_{x} 90^{\circ} y 0^{\circ}{ }_{y} 90^{\circ}{ }_{x}\right)$. Relaxation of longitudinal two-spin order is detected if the sequence between points A and B in (a) is replaced with (b); the proton composite pulse is $\left(90^{\circ}{ }_{x} 90^{\circ} y 90^{\circ} y 90^{\circ}\right.$ ). Composite pulses are applied at the center of amide ${ }^{1} \mathrm{H}$. Narrow and wide solid bars are $90^{\circ}$ and $180^{\circ}$ pulses, respectively. All pulses have $x$ phase unless otherwise stated. Phase cycles are $\phi_{1}=x,-x ; \phi_{2}=x, x,-x,-x ; \phi_{3}=$ $\mathrm{y} ; \phi_{4}=x$; and receiver phase $=-x, x, x,-x$. Gradients G4 and G7 are used for coherence selection; other gradients are for artifact suppression. Echoantiecho quadrature detection is achieved by inverting $\phi_{3}, \phi_{4}$, and the sign of gradient G4; ${ }^{12}$ axial peaks are shifted by inverting $\phi_{1}$, and the receiver phase. Other delays are $\Delta=2.7 \mathrm{~ms}, \xi=950 \mu \mathrm{~s}, \epsilon=2.675 \mathrm{~ms}, \zeta=250$ $\mu \mathrm{s}$.

Pulse sequences for measuring $R_{2}{ }^{\alpha}, \eta_{x y}$, and $R_{1}{ }^{2 \mathrm{HzNz}}$ in large proteins are shown in Figure 1. During the Hahn echo period of length $2 \tau$ in Figure 1a, the narrow and broad doublet components relax independently because cross relaxation between the doublet components is suppressed when $2 \tau=N / J_{\mathrm{NH}}$ and $2 \pi J_{\mathrm{NH}} \gg R_{1}{ }^{\mathrm{H}} / 2$, in which $N$ is an integer and $J_{\mathrm{NH}}$ is the scalar coupling constant. ${ }^{11}$ The $\alpha$ and $\beta$ spin states of ${ }^{1} \mathrm{H}$ spins are interchanged by a $180^{\circ}$ proton pulse. Therefore, if the proton composite pulse element in Figure 1a is effectively a $0^{\circ}$ or $180^{\circ}$ pulse (see Figure 1 caption for details), the intensity of the narrow or broad doublet component is recorded by the TROSY scheme following point $B$, respectively. Figure 1 b shows modifications of Figure 1a necessary to measure $\mathrm{R}_{1}{ }^{2 \mathrm{HzNz}}$ during a relaxation period $\tau$.

The intensity decays for the three experiments are given by:

$$
\begin{align*}
I^{\alpha}(2 \tau) & =I(0) \exp \left[-2 \tau R_{2}^{\alpha}\right]  \tag{3}\\
I^{\beta}(2 \tau) & =I(0) \exp \left[-2 \tau R_{2}^{\beta}\right]  \tag{4}\\
I^{\mathrm{HzN} z}(\tau) & =I(0) \exp \left[-\tau R_{1}^{2 \mathrm{HzN} \mathrm{Z}}\right] \tag{5}
\end{align*}
$$

in which the prefactor $I(0)$ is the same in eqs $3-5$ because the three pulse sequences have identical delays and number of pulses when $\tau=0$. Thus, the rate constants needed to determine $R_{\mathrm{ex}}$ using


Figure 2. Measurements of (a) $\eta_{x y}$ and (b) $R_{\mathrm{ex}}$ for $\left[85 \%-\mathrm{D}, \mathrm{U}-{ }^{15} \mathrm{~N}\right] E$. coli RNaseH using the pulse sequences in Figure 1. Values of $\eta_{x y}$ obtained from the pulse sequence depicted in Figure 1a, denoted as $\eta_{x y}{ }^{\mathrm{T}}$, are plotted versus values measured from conventional sequences, denoted as $\eta_{x y}{ }^{\mathrm{C}}$. The inset of Figure 2b compares the $R_{\text {ex }}$ obtained from conventional $R_{2}$ and $\eta_{x y}$ measurements ( $R_{\mathrm{ex}}{ }^{\mathrm{c}}$ ) and from the new sequences ( $R_{\mathrm{ex}}{ }^{\mathrm{T}}$ ). The relaxation delay $2 \tau=10 / J_{\mathrm{NH}}=108 \mathrm{~ms}$, corresponding to $I^{\beta}(2 \tau) / I^{\alpha}(2 \tau)=0.3$ for most residues. $\kappa$ was determined by the trimmed mean of $1+\left(R_{2}{ }^{\alpha}-R_{1}{ }^{2 \mathrm{HzNz}} / 2\right.$ $\left.+R_{1} \mathrm{~N} / 2\right) / \eta_{x y}$ to be 1.4. For nonexchanging residues, $R_{\text {ex }}$ is distributed around zero with a standard deviation of $0.6 \mathrm{~s}^{-1}$ due to variation in magnitude and orientation of the ${ }^{15} \mathrm{~N}$ CSA tensor. ${ }^{8}$ Data were collected in triplicate with a 1.0 mM sample at 310 K on a Bruker DRX600 spectrometer equipped with a triple resonance probe.
eq 2 are obtained from the intensity ratios:

$$
\begin{gather*}
I^{\beta}(2 \tau) / I^{\alpha}(2 \tau)=\exp \left[-4 \tau \eta_{x y}\right]  \tag{6}\\
I^{\alpha}(2 \tau) / \mathrm{I}^{2 \mathrm{HzNz}}(\tau)=\exp \left[-2 \tau\left(R_{2}^{\alpha}-R_{1}^{2 \mathrm{HzNz}} / 2\right)\right] \tag{7}
\end{gather*}
$$

The rate constants are determined by curve-fitting the intensity ratios measured at different values of $\tau$, or by repeatedly measuring duplicate spectra with a single value of $\tau .{ }^{13}$

The new experiments incorporate a number of features to improve sensitivity and accuracy for larger proteins. First, TROSY (transverse relaxation optimized spectroscopy) is employed during the $t_{1}$ evolution and $t_{2}$ detection periods for improved resolution and sensitivity in large proteins. ${ }^{5} R_{2}{ }^{\alpha}$ and $R_{2}{ }^{\beta}$ are measured independently to obtain $\eta_{x y}$, because independent detection of inphase and anti-phase ${ }^{15} \mathrm{~N}$ magnetization, as used in conventional experiments for measuring $\eta_{x y},{ }^{7,14}$ is difficult to utilize with TROSY. Second, ${ }^{1} \mathrm{H} 180^{\circ}$ pulses or composite pulse sequences are not used for decoupling during the Hahn echo period. Decoupling can interfere with echo formation and increase apparent relaxation rate constants. ${ }^{11}$ Third, data are not recorded for $\tau=0$; therefore, a minimum of three 2D spectra are sufficient to determine $R_{\text {ex }}$ using eqs 2,6 , and 7 . Fourth, the uncertainty in $R_{\text {ex }}$ contributed by the $\eta_{x y}$ measurement is scaled down by a factor of $\kappa /(\kappa-1)$, which is more than 6-fold at 800 MHz , compared with approaches that obtain $R_{\text {ex }}=R_{2}-\kappa \eta_{x y}$ from $R_{2}$ and $\eta_{x y .}{ }^{9}$ As in all experiments involving anti-phase ${ }^{15} \mathrm{~N}$ magnetization, the sensitivity of these sequences may be reduced for solvent-exposed residues by rapid amide proton exchange, especially at high pH .

The proposed technique was validated using RNaseH. As shown in Figure 2, good agreement is obtained between $\eta_{x y}$ and $R_{\text {ex }}$ obtained from the new and conventional pulse sequences. ${ }^{7}$ Chemical exchange is detected for residues Lys60 and Trp90, in agreement with previous reports. ${ }^{7}$ Small exchange contributions are also evident for residues near positions 10 and 120.

The new methods were applied to glycerol 3-phosphate (G3P) bound TIM to illustrate the advantages for larger proteins. In TIM, the average ${ }^{15} \mathrm{~N} R_{2}$ is $37 \mathrm{~s}^{-1}$ at 800 MHz at 298 K , while the average $R_{2}{ }^{\alpha}$ is only $6 \mathrm{~s}^{-1}$. Therefore, the relative contribution of $R_{\mathrm{ex}}$ to $R_{2}{ }^{\alpha}$ can be 6 times as high as the contribution to $R_{2}$. As shown in Figure 3 , chemical exchange broadening caused by the well-known loop


Figure 3. Chemical exchange in G3P bound [ $85 \%$-D, U- $\left.{ }^{15} \mathrm{~N}\right]$ TIM detected using the pulse sequences in Figure 1. The insert is an expanded view of the region encompassing residues in the active site loop 6. The relaxation delay $2 \tau=2 / J_{\mathrm{NH}}=21.6 \mathrm{~ms}$, corresponding to $I^{\beta}(2 \tau) / I^{\alpha}(2 \tau)=0.3$ for most residues. $\kappa$ was determined by the trimmed mean of $1+\left(R_{2}{ }^{\alpha}-R_{1}{ }^{2 \mathrm{HzNz}} / 2\right) /$ $\eta_{x y}$ to be 1.2. For nonexchanging residues, $R_{\text {ex }}$ is distributed around zero with a standard deviation of $1.7 \mathrm{~s}^{-1}$ due to variation in magnitude and orientation of the ${ }^{15} \mathrm{~N}$ CSA tensor. ${ }^{8}$ The data were collected in quadruplicate on a 1 mM TIM (with 10 mM G3P) sample at 800 MHz using a Bruker DRX800 NMR spectrometer equipped with a triple-resonance cryoprobe.
motion in TIM ${ }^{15}$ is observed for residues at the N (residues V167 and W168) and C termini (residues L174) hinge of loop 6. Other residues (T177, N213, F220, and K221) with significant chemical exchange are located in the vicinity of loop 6 in the 3D structure of TIM. These results are consistent with the hypothesis that the termini of the loop act as the hinges. ${ }^{15}$

The techniques presented herein allow detection of conformational exchange processes in proteins with MW $>50 \mathrm{kD}$. Thus, these methods, along with the TROSY-CPMG experiment, ${ }^{16,17}$ expand the capacity of NMR spectroscopy to investigate $\mu \mathrm{s}-\mathrm{ms}$ time scale motions involved in protein function.

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